

Heavy-light hadrons: Old laces & new pieces

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“Solomon saith: There is no new thing upon the earth. So that as Plato had an imagination, that all knowledge was but remembrance; so Solomon giveth his sentence, **that all novelty is but oblivion.**

Francis Bacon: Essays, LVIII”

— Jorge Luis Borges, The Aleph and Other Stories

Outline:

- **Old Laces** - Chiral dynamics, summarized in e.g. "Nuclear Chiral Dynamics"
(M. Rho, I. Zahed, MAN; Singapore 1996)
- **Two revolutions** (exp. and theoretical)
- **New Pieces** - Holographic HL Chiral Effective Action
Y. Liu, I. Zahed; PRD 95(2017) 056022, PRD 95(2017) 116012
Y. Liu, MAN, I. Zahed; PRD 100(2019) 126023; PRD .. (2022) in press (Tetrequarks)
Y. Liu, MAN, I. Zahed; PRD 104(2021), 114021; PRD 104(2021) 114022 (Pentaguerks)
Y. Liu, K. Mamo, MAN, I. Zahed; PRD 104(2021) 114023. (Photoproduction of Pentaguerks)

Heavy and light.

2 (3) quarks (u, d, s)

3 heavy quarks (c, b, t)

$$m_{u,d} \ll \Lambda_{QCD}, \quad m_s \sim \Lambda_{QCD}$$

$$m_h \gg \Lambda_{QCD}$$

• Light sector: SBXS $SU_L(n_L) \times SU_R(n_R) \rightarrow SU_V(n_c)$

Effective chiral actions (G-model, NJL, X.P.T.,)

• Heavy sector: heavy spin decouples as $m_h \rightarrow \infty$

Then $(0^-, 1^-)$ becomes degenerated in the $m_h \rightarrow \infty$ limit

$$H = \frac{1 + \gamma_5}{2} (i \gamma_5 D + \gamma^\mu D_\mu^*)$$

• Expansion in $1/N_c$

Light sector

- Goldstone particles (pions)

$$\mathcal{L}_G \sim \frac{f_\pi^2}{4} \text{Tr} \partial_\mu U \partial^\mu U + \dots$$

$$U = e^{i/f_\pi \vec{c} \cdot \vec{\pi}}$$

- What about baryons? "Ergonomic solution": Skyrme
Baryon is a soliton, Baryon number is topological (winding)

$$\vec{\pi} = f(r) \frac{\vec{r}}{r}, \quad \mathcal{L} = \mathcal{L}_G + \mathcal{L}_{\text{Skyrme}} + \mathcal{L}_{\text{WZW}} \xrightarrow{(\pi^0 \rightarrow \gamma\gamma)} \text{stabilizer}$$

- Such classical solution

has zero modes (rotations $U \rightarrow A(t) U A^\dagger(t)$)

- Quantization of collective coordinates in moduli space (Witten, Adkins, Napp)

$$H = M_0 + \frac{I^2}{2\Omega_{\text{sol}}} = M_0 + \frac{J^2}{2\Omega_{\text{sol}}}$$

- Two parameters, good phenomenology

Strange sector

- $SU(2) \rightarrow SU(3)$: Mazur, MAN, Presniewicz $\rightarrow$ "elegant disaster" (1984)

- Way-out: m_s is too heavy:

BORN-Oppenheimer approx: fast vibrations of kaon in $SU(2)$ solitonic background. Then slow rotation of the bound state happens in presence of Berry phase from kaon

$$H = M_0 + \text{binding} + \frac{(\bar{J} - (1 - c_K) \text{tr } K \bar{I} K^\dagger)^2}{2 \Omega_{\text{sol}}}$$

(isospin-spin transmutation)



versus



Callan-Klebanov - (Hornbostel) 1985

- Successful phenomenology (also, no pentaquark)

Heavy sector:

- One has to remember about heavy spin symmetry

$$H = M_0 + \text{binding} + \frac{[(\vec{J} - \vec{S}_H) - (1 - c_D) \text{tr } \tilde{D} \tilde{I} D^\dagger - (1 - c_D^*) \text{tr } D^* \tilde{I} D^{*\dagger}]^2}{2 \Omega_{\text{sol}}}$$

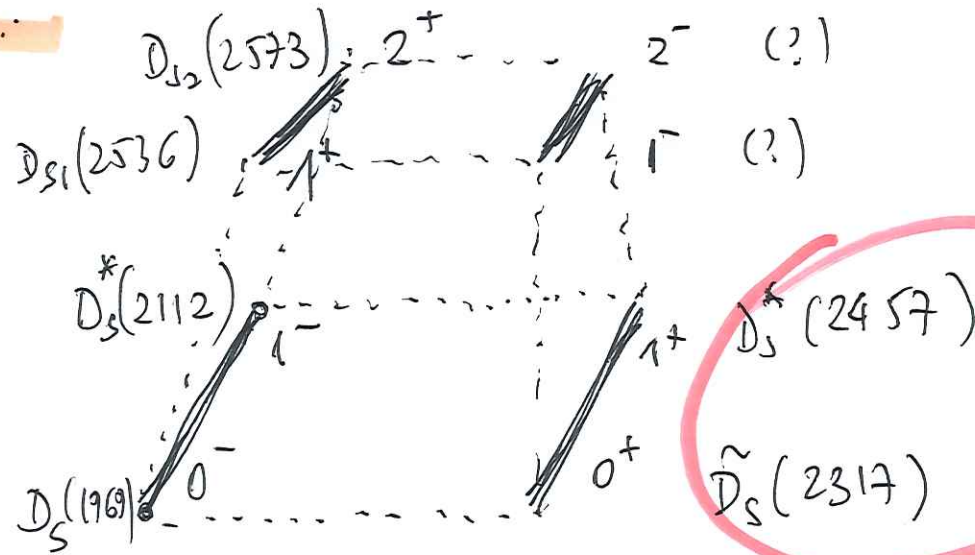
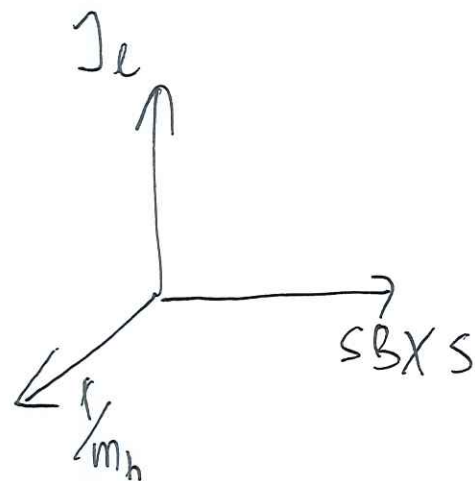
- When $m_h \rightarrow \infty$, phases cancel $\frac{(\vec{J} - \vec{S}_H)^2}{2 \Omega_{\text{sol}}} = \frac{I^2}{2 \Omega_{\text{sol}}}$

(Heavy spin symmetry at the baryonic level)

- Soliton can capture more than one meson (notice), Savage-Wise symmetry
- Same happens (e priori) for χ doubler

Chiral doublers scenario (MAN, M.Rho, I.Zahed (1992), Bardeen-Hill (1993))

- Requirement of simultaneous presence of both symmetries enforces the emergence of opposite to tl , multiplet $G = \frac{1+\gamma}{2} (\tilde{D} + \gamma^\mu \gamma_5 \tilde{D}_\mu^*)$
- Doublers communicate via axial current
- $m_G - m_H \sim O(\Sigma_{light}) \sim 350 \text{ MeV}$
- Reorganization, not doubling the hadronic states
- Spectacular example:



Beber

Meson-Baryon "supersymmetry"

- QQ  $\longrightarrow$ Doublet with color $\bar{3}$, like $\bar{Q}$
- Such symmetry relates spin-splittings of $\bar{H}L$ to HHL
($\Delta M_{HHL} = \frac{3}{4} \Delta M_{\bar{H}L}$) Sonnerup & Wise (1990)
- Interesting duality $HH \leftrightarrow \bar{H}$ leads to further correspondence
e.g. $HLL \rightarrow \bar{H}HLL$ (heavy baryon - doubly heavy tetras.)

New Zoo (alike quark chemistry)

- soliton captures H meson
- soliton — " — G meson
- soliton captures $\bar{H}$ (heavy pentaquark)
- soliton — " — $\bar{G}$ (darken of h. pentaquark)
- soliton captures more than one
- etc ...

Exotica

$\bar{c}du\bar{s}$
 $c\bar{c}q\bar{q}$
 $c\bar{c}u\bar{d}$
 $c\bar{c}u\bar{s}$
 $c\bar{c}s\bar{s}$
 $c\bar{c}c\bar{c}$
 $b\bar{b}u\bar{d}$
 $c\bar{c}uud$
 $c\bar{c}uds$
 $ccu\bar{d}$

States

2866 2909

$X_0(2900), X_1(2900) [21, 22]$
 $\rightarrow X_{c1}(3872) [6]$

$\left\{ \begin{array}{l} 2007 \quad 1864.7 \\ D^{*0} \quad \bar{D}^0 \end{array} \right.$ but wider threshold of (120 keV)

$\left\{ \begin{array}{l} Z_c(3900) [23], Z_c(4020) [24, 25], Z_c(4050) [26], X(4100) [27], \\ Z_c(4200) [28], Z_c(4430) [29-32], R_{c0}(4240) [31] \end{array} \right.$

$Z_{cs}(3985) [33], Z_{cs}(4000), Z_{cs}(4220) [34]$

$\left\{ \begin{array}{l} X_{c1}(4140) [35-38], X_{c1}(4274), X_{c0}(4500), X_{c0}(4700) [38], \\ X(4630), X(4685) [34], X(4740) [39] \end{array} \right.$

$X(6900) [14]$

$Z_b(10610), Z_b(10650) [40]$

$P_c(4312) [41], P_c(4380) [42], P_c(4440), P_c(4457) [41],$

$P_c(4357) [43] \quad 36 \quad \text{weekend}$

$P_{cs}(4459) [44]$

T_{cc}^+

$\begin{array}{c} DD^{*+} \\ \bar{c} \quad \bar{c}d \\ cu \quad 1 \\ 1865 \quad 2010 \\ 3875 \end{array}$

● Gravity / Gauge duality (holography, AdS/CFT)

In the 70'ies:

QCD
Fundamental theory of quarks & gluons

strings - flux tubes


effective, like LUND model

Maldacene

Gauge theory in 4 dim $\simeq$ String theory in higher dimensions

Witten:

Application of duality to QCD - pure YM in 3+1
at large N and $\lambda = g_{YM}^2 N$

Surprising similarity of spectrum of glueballs to large N lattice

Sakai & Sugimoto 2005

- Adding N_f massless fermions: $SU_L(N_f) \times SU_R(N_f) \rightarrow SU_V(N_f)$ $\Delta B \neq 0$

- Low energy limit $S = S_{YM} + S_{CS}$

$$S_{YM} = \kappa \int d^4x dz \operatorname{Tr} \left(\frac{1}{2} k(z)^{-1/3} F_{\mu\nu}^2 + k(z) F_{yz}^2 \right) \quad k(z) = 1 + z^2$$

Mode expansion:

$$A_\mu(x^\mu, z) = \sum_n B_\mu^{(n)}(x^\mu) \Psi_n(z)$$

$$A_5(x^\mu, z) = \sum_n f^{(n)}(x^\mu) \Phi_n(z).$$

- Keeping only $f^{(0)}$ gives $\mathcal{L} = \mathcal{L}_S + \mathcal{L}_{\text{Skyrme}} + \mathcal{L}_{W2}$, i.e. Skyrme model
 $B_\mu^{(1)} \sim g$, $B_\mu^{(2)} \sim a_1$ give hidden gauge model

- Successful phenomenology with very few parameters.

From S. Sugimoto talk

mass

mass	ρ	a_1	ρ'	(a'_1)	ρ''
exp.(MeV)	776	1230	1465	(1640)	1720
our model	[776]	1189	1607	2023	2435
ratio	[1]	1.03	0.911	(0.811)	0.706

↑
input ($M_{KK} \simeq 949$ MeV)

coupling

coupling		fitting m_ρ and f_π	experiment
f_π	$1.13 \cdot \kappa^{1/2} M_{KK}$	[92.4 MeV]	92.4 MeV
L_1	$0.0785 \cdot \kappa$	0.584×10^{-3}	$(0.1 \sim 0.7) \times 10^{-3}$
L_2	$0.157 \cdot \kappa$	1.17×10^{-3}	$(1.1 \sim 1.7) \times 10^{-3}$
L_3	$-0.471 \cdot \kappa$	-3.51×10^{-3}	$-(2.4 \sim 4.6) \times 10^{-3}$
L_9	$1.17 \cdot \kappa$	8.74×10^{-3}	$(6.2 \sim 7.6) \times 10^{-3}$
L_{10}	$-1.17 \cdot \kappa$	-8.74×10^{-3}	$-(4.8 \sim 6.3) \times 10^{-3}$
$g_{\rho\pi\pi}$	$0.415 \cdot \kappa^{-1/2}$	4.81	5.99
g_ρ	$2.11 \cdot \kappa^{1/2} M_{KK}^2$	0.164 GeV^2	0.121 GeV^2
$g_{a_1\rho\pi}$	$0.421 \cdot \kappa^{-1/2} M_{KK}$	4.63 GeV	$2.8 \sim 4.2 \text{ GeV}$

Baryon in Sakai & Sugimoto

4 dim pion Skyrmion (static solution)

5 dim gauge field $\downarrow$ BPS T instanton in x_1, x_2, x_3, z in flavor

Topological number $\equiv$ baryon number

Direct realization of 1989 Atiyah-Manton idea

$$U(\vec{x}) = P \exp \left(i \int dz A_z(\vec{x}, z) \right)$$

8 zero modes $(SU(2), \vec{x}, z, \theta)$ lead to moduli space

quantization

$$M = M_0 + \left(\sqrt{\frac{(l+1)^2}{6} + \frac{2}{15} N_c^2} + \sqrt{\frac{2}{3} (n_S + n_Z + 1)} \right) M_{KK}$$

$$l = 2I = 2J = 1, 3, 5, \dots$$

From S. Sugimoto talk

- If we choose $M_{KK} \simeq 500$ MeV
and use nucleon mass ($\simeq 940$ MeV) to fix the constant M_0 ,
(we only consider the mass difference), we obtain

(n_ρ, n_z)	(0,0)	(1,0)	(0,1)	(1,1)	(2,0)/(0,2)	(2,1)/(0,3)	(1,2)/(3,0)
$N(l=1)$	[940] ⁺	1348 ⁺	1348 ⁻	1756 ⁻	1756 ⁺ , 1756 ⁺	2164 ⁻ , 2164 ⁻	2164 ⁺ , 2164 ⁺
$\Delta(l=3)$	1240 ⁺	1648 ⁺	1648 ⁻	2056 ⁻	2056 ⁺ , 2056 ⁺	2464 ⁻ , 2464 ⁻	2464 ⁺ , 2464 ⁺

↑
States appeared in the Skyrme model (± : parity)

- $I = J$ states from Particle Data Group look like....

(n_ρ, n_z)	(0,0)	(1,0)	(0,1)	(1,1)	(2,0)/(0,2)	(2,1)/(0,3)	(1,2)/(3,0)
$N(l=1)$	940 ⁺	1440 ⁺	1535 ⁻	1655 ⁻	1710 ⁺ , ?	2090 _* ⁻ , ?	2100 _* ⁺ , ?
$\Delta(l=3)$	1232 ⁺	1600 ⁺	1700 ⁻	1940 _* ⁻	1920 ⁺ , ?	?, ?	?, ?

(? : not found, * : evidence of existence is poor)

S. Segimoto talk (Moniaul)

Table 2: $\langle r^2 \rangle_{I=0}$, $\langle r^2 \rangle_{I=1}$ and $\langle r^2 \rangle_A$ are isoscalar, isovector and axial mean square radii, respectively. $g_{I=0}$ and $g_{I=1}$ are isoscalar and isovector g-factors. g_A is the axial coupling.

	our model	experiment
$\langle r^2 \rangle_{I=0}$	$(0.74 \text{ fm})^2$	$(0.81 \text{ fm})^2$
$\langle r^2 \rangle_{I=1}$	$(0.74 \text{ fm})^2$	$(0.94 \text{ fm})^2$
$\langle r^2 \rangle_A$	$(0.54 \text{ fm})^2$	$(0.67 \text{ fm})^2$
$g_{I=0}$	1.7	1.8
$g_{I=1}$	7.0	9.4
g_A	0.73	1.3

Holographic H-L effective theory (Liu, Zahed; 2017)

- "CK"-like scheme based on adding ∞ -heavy flavor ($m_H = \infty$)

$$M = M_0 + (N_Q + N_{\bar{Q}}) \underline{m_H} + \left[\frac{(l+1)^2}{6} + \frac{2}{15} N_c^2 \left(1 - \frac{15(N_Q - N_{\bar{Q}})}{4 N_c} + \frac{5(N_Q - N_{\bar{Q}})^2}{3 N_c^2} \right) \right]^{\frac{1}{2}} M_{KK} + \sqrt{\frac{2}{3}} (n_S + n_Z + 1) M_{KK}$$

- Various combinations of q -numbers give all types of H-L holograms.

$N_Q = 1$	$N_{\bar{Q}} = 0$	baryons	HLL
$N_Q = 2$	$N_{\bar{Q}} = 0$	- - -	HHL
$N_Q = N_{\bar{Q}} = 1$		Pentagons	H $\bar{H}$ LLL
$n_Z \neq 0$	excited	$n_S \neq 0$	odd parity ...

- Three parameters: M_{KK} , M_0 , m_H

Adding spin-effects (subleading terms in m_H^{-1}) Liu, MAN, Zahed (2021)

- 3 parameters $M_0 \rightarrow m_N$, $M_{KK} \rightarrow m_{\Lambda_c}$, $m_H = M_D$ for c
 $m_H = M_B$ for b

- 3 Pentaquarks $\frac{1}{2} \frac{1}{2}^-$, $\frac{1}{2} \frac{1}{2}^-$, $\frac{1}{2} \frac{3}{2}^-$ $I \supset \pi$
 $s=1$ $s=0$ $s=1$

($\frac{1}{2} \frac{5}{2}^\pm$ ruled out), consistent with $P_c(4312, 4440, 4457)$ (LHCb)

- Recently reported $P_c(4337)$ at 3 σ significance is not supported

- Open and hidden decay widths Liu, MAN, Zahed (2021) b

e.g. $P_c \rightarrow \Lambda_c + \bar{D}$ $\Gamma(s=0, J=\frac{1}{2}) : \Gamma(s=1, J=\frac{1}{2}) : \Gamma(s=1, J=\frac{3}{2}) = \frac{1}{2} : \frac{5}{6} : \frac{1}{3}$

- Formfactors (Liu, Memo, MAN, Zahed; 2022), consistent with recent GLUEX result.
 $\gamma p \rightarrow \underbrace{J/\psi p}_{P_c^+}$

TABLE I. Charm baryons and Pentaquarks

B	IJ^P	l	n_ρ	n_z	N_Q	$N_{\bar{Q}}$	Mass-MeV	Exp-MeV
Λ_c	$0\frac{1}{2}^+$	0	0	0	1	0	2286	2286
Σ_c	$1\frac{1}{2}^+$	2	0	0	1	0	2557	2453
	$1\frac{3}{2}^+$	2	0	0	1	0	2596	2520
Λ_c^*	$0\frac{1}{2}^-$	0	0	1	1	0	2683	2595
	$0\frac{1}{2}^+$	0	1	0	1	0	2726	2765
Σ_c^*	$1\frac{1}{2}^-, 1\frac{3}{2}^-$	2	0	1	1	0	[2947/2986]	—
	$1\frac{1}{2}^+, 1\frac{3}{2}^+$	2	1	0	1	0	[2948/2995]	—
P_c	$\frac{1}{2}\frac{1}{2}^-, \frac{1}{2}\frac{3}{2}^-$	1	0	0	1	1	[4340/4360/4374]	[4312/4440/4457]
P_c^*	$1\frac{1}{2}^-, 1\frac{3}{2}^-$	1	0	1	1	1	[4732/4752/4767]	—
	$1\frac{1}{2}^+, 1\frac{3}{2}^+$	1	1	0	1	1	[4725/4746/4763]	—

To fix the parameters for the charmed heavy baryons, we choose $M_D = 1.87$ GeV for the D-meson mass in (7) and fix $M_{KK} = 0.475$ GeV to reproduce the $M_{\Lambda_c} = 2.286$ GeV. This low value of M_{KK} is consistent with the value used to reproduce the nucleon

TABLE II. Bottom baryons and Pentaquarks

B	IJ^P	l	n_ρ	n_z	N_Q	$N_{\bar{Q}}$	Mass-MeV	Exp-MeV
Λ_b	$0\frac{1}{2}^+$	0	0	0	1	0	5608	5620
Σ_b	$1\frac{1}{2}^+$	2	0	0	1	0	5962	5810
	$1\frac{3}{2}^+$	2	0	0	1	0	5978	5830
Λ_b^*	$0\frac{1}{2}^-$	0	0	1	1	0	5998	5912
	$0\frac{1}{2}^+$	0	1	0	1	0	6029	(6072)
Σ_b^*	$1\frac{1}{2}^-, 1\frac{3}{2}^-$	2	0	1	1	0	[6351/6367]	—
	$1\frac{1}{2}^+, 1\frac{3}{2}^+$	2	1	0	1	0	[6344/6367]	—
P_b	$\frac{1}{2}\frac{1}{2}^-, \frac{1}{2}\frac{3}{2}^-$	1	0	0	1	1	[11155/11163/11167]	—
P_b^*	$1\frac{1}{2}^-, 1\frac{3}{2}^-$	1	0	1	1	1	[11544/11553/11556]	—
	$1\frac{1}{2}^+, 1\frac{3}{2}^+$	1	1	0	1	1	[11532 /11543/11579]	—

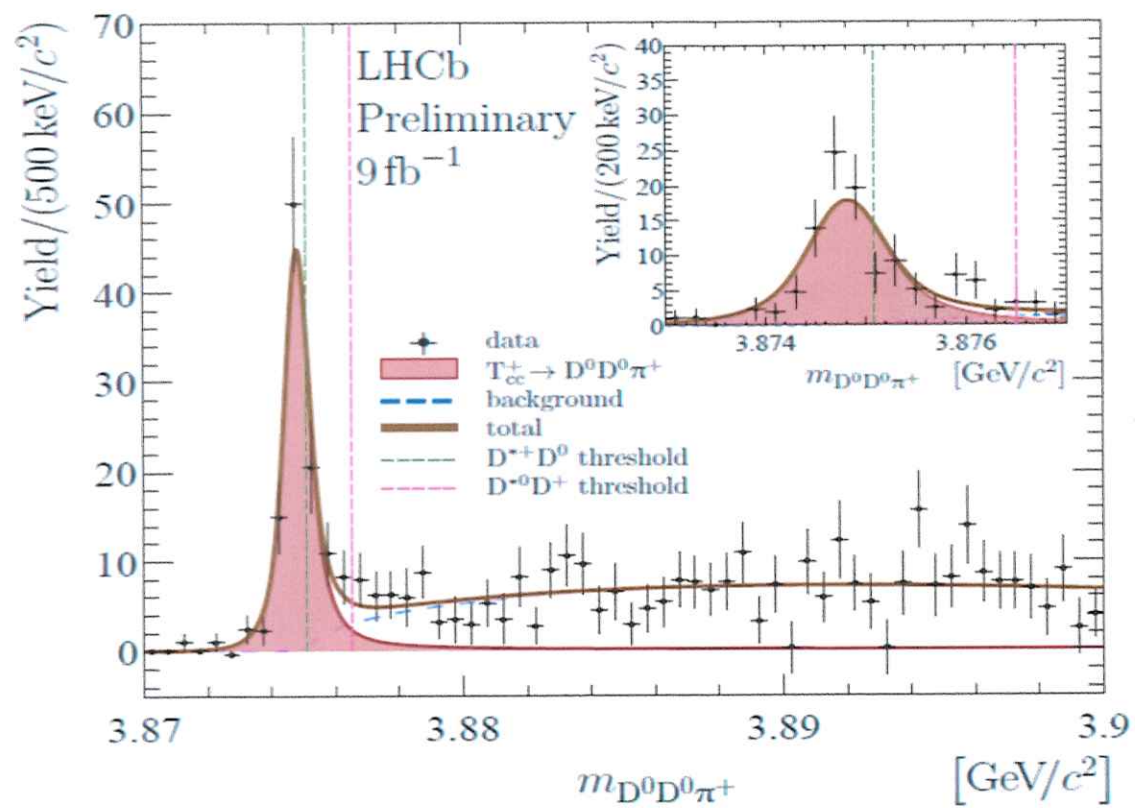


FIG. 1. Newly measured isoscalar charmed T_{cc}^+ tetraquark by LHCb [48].

Tetraqark puzzle

- $\overline{HHLL}$ — several predictions (positive/negative within $\pm 200 \text{ MeV}$)
- Measurement of $\Xi_{cc}^{++} (3621)$ fixed the normalization for bound tetraquarks.
(Kerlin, Rosner 2017, Eichten-Quigg (2017)), bounded up to 200 MeV (!).
- In holographic picture, role of the instanton (baryons, pentaquarks) is played by sphaleron (Baryon # = 0), binding two heavy mesons.
(Liu, MAN, Zahed; 2019) $b\bar{b}q\bar{q}$ bounded 80 MeV, $ccq\bar{q}$ bounded by 90 MeV
- T_{cc}^+ (01^+), narrow, discovered at LHCb, bounded at -360 keV , $\Gamma \sim 50 \text{ keV}$
- Normalizing mass to T_{cc}^+ , we (Liu, MAN, Zahed, 2022) predict mass of T_{bc} and T_{bb} , and calculate very narrow width. ($m_H = \infty$)

TABLE I. Binding energies for tetraquarks versus the 't Hooft coupling $\lambda = g_{\text{YM}}^2 N_c$ with $M_{\text{KK}} = 1$ GeV

λ	$QQ\bar{q}\bar{q}$ GeV	$bb\bar{q}\bar{q}$ GeV	$bc\bar{q}\bar{q}$ GeV	$cc\bar{q}\bar{q}$ GeV
10	-0.097	-0.088	-0.080	-0.072
15	-0.107	-0.091	-0.077	-0.062
20	-0.108	-0.085	-0.064	-0.041
25	-0.103	-0.073	-0.045	-0.018
30	-0.093	-0.056	-0.024	-0.0016
32	-0.089	-0.048	-0.015	0.00073

TABLE II. Binding energies for tetraquarks versus the 't Hooft coupling $\lambda = g_{\text{YM}}^2 N_c$ with $M_{\text{KK}} = 0.475$ GeV.

λ	$QQ\bar{q}\bar{q}$ GeV	$bb\bar{q}\bar{q}$ GeV	$bc\bar{q}\bar{q}$ GeV	$cc\bar{q}\bar{q}$ GeV
10	-0.046	-0.044	-0.042	-0.040
15	-0.051	-0.047	-0.044	-0.040
20	-0.051	-0.046	-0.040	-0.035
25	-0.049	-0.042	-0.035	-0.028
30	-0.045	-0.035	-0.027	-0.019
40	-0.031	-0.018	-0.0076	0.0011

Summary:

- Strongly coupled QCD could be approached via duality from string theory in large N_c and large λ limit, including spectra of heavy-light hadrons
- Few parameters and very restrictive predictions, so models are confrontable
- Approach based on confinement, SBXS and heavy spin symmetry, in the limit of large N_c and large 't Hooft coupling.
- "High brow" theory boils down to relatively simple QM in moduli space.
- Astonishing and deep analogies to "old" physics.

particles

solar - sugimoto

table

our table (I)

