Rapidity space entanglement and high energy process

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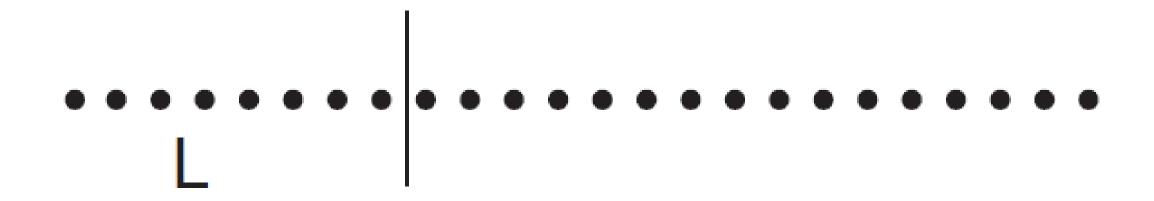
• The talk is based on recent work 2203.00739 and 2202.02612.

Outline

- Introduction to Quantum Entanglement
- High-energy process: large rapidity limit
- Rapidity space entanglement: sub-critical system
- Rapidity space entanglement : critical system and rapidity evolution
- Rapidity space entanglement: string picture
- Conclusion

Introduction to Quantum Entanglement

- $\mathcal{H} = \mathcal{H}_0 \otimes \mathcal{H}_1$, full density matrix ρ .
- Reduced density matrix : $\rho_0 = tr_1 \rho$.
- Entanglement entropy: $S = -tr\rho_0 \ln \rho_0$. Roughly, about
 - 1. How measurement in 0 affects 1.
 - 2. How generic the state is.
 - 3. How many channels between subsystems.
- Reflects deep laws in quantum many body system, in particular, QFT.



Introduction to Quantum entanglement • The real space entanglement focuses on a segment in position space with length L.

Introduction to Quantum Entanglement

- Example: real space entanglement in 1D.
- 1. Generic state: $S = O(L) = \ln \dim \mathcal{H}$. All 2^L sates contribute.
- 2. Short range interaction: Area law, $S \leq O(\ln L)$.
- Gapless system, $c(L) = L \frac{dS}{dL} = c + O(\frac{1}{L})$.
- Gapped system, $c(L) = L \frac{dS}{dL} = O(e^{-mL})$ at large distance.

Introduction to Quantum Entanglement

- Example: real space entanglement in 1D.
- 3. Non-trivial even for free-system.
- Macroscopic: replica trick, CFT based methods.
- Microscopic: Fisher-Hartwig conjecture.
 Open questions remain.
- 4. Relation to matrix-product states, entanglement renormalization group and string duality.

High-energy process: large rapidity limit

- 1. High-energy experiment: presence of a large rapidity gap *Y*.
- DIS at small- $x: Y = \ln \frac{1}{x}$.
- Forward scattering with $s \gg -t$, $Y = \ln \frac{s}{m^2}$.
- 3. Very non-trivial asymptotic behavior in *Y*.
- Perturbative evolution equations and saturation conjecture.
- Insights from String-Gauge duality.

High-energy process: large rapidity limit

- 4. Nontrivial "conspiracy" between fast and slow degrees freedom.
- 5. How many entanglement between fast and slow degrees of freedom?

Rapidity space entanglement: sub-critical system

- Meson state in the 2D QCD in large N_c limit.
- 1. Light-front wave function: $|m\rangle = \frac{1}{\sqrt{N}} \sum \phi(x) |x, \bar{x}\rangle$.
- 0 < x < 1: longitudinal momentum fraction.
- *N*: total number of digits.
- 2. Rapidity space entanglement between sub-systems: $[0,1]N = [0,x_0]N \cup [x_0,1]N$.

x_0N $(1-x_0)N$

Rapidity space entanglement : subcritical system • In the light-front formulation of QFT, there are N total digits in longitudinal momentum fraction. We consider entanglement between the first x_0N and the rest.

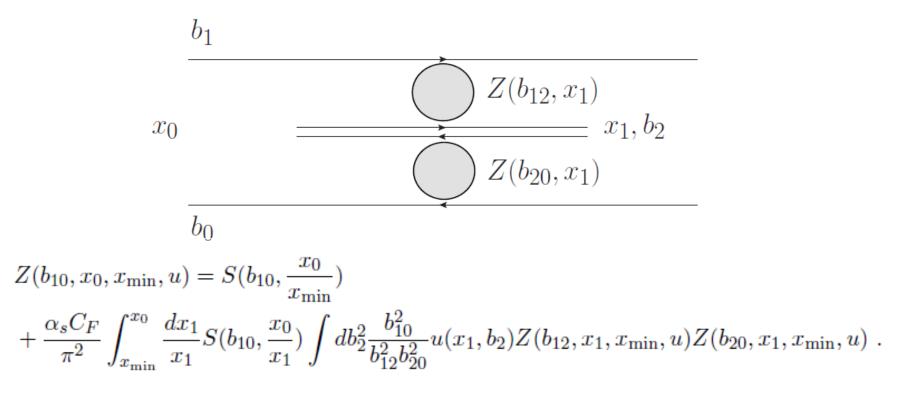
Rapidity space entanglement: sub-critical system

- 3. $S(x_0) = c(x_0) \ln x_0 N + b(x_0)$. Area law satisfied.
- 4. $c(x_0)$ and $b(x_0)$ expressed in terms of quark parton distribution functions (PDFs).
- 5. For small x_0 , $c(x_0) \sim x_0^{2\beta+1}$ and $b(x_0) \sim x_0^{2\beta+1} \ln \frac{1}{x_0}$
- The same asymptotic coefficients for forward scattering: $A(s) \sim s^{-2\beta}$.

Rapidity space entanglement: sub-critical system

- $S(x_0) = c(x_0) \ln x_0 N + b(x_0)$.
- 1. The "c function": $c(x_0) = \int_0^{x_0} q(x) + \overline{q}(x) dx \le 1$.
- 2. Expected to generalized to multi-parton wave functions in sub-critical system.

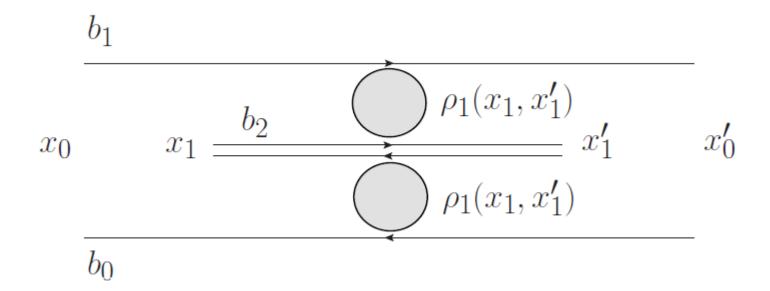
- In QCD, a famous example is the soft gluon wave function of a quarkonium system.
- 1. Emission of small-*x* gluon generates rapidity divergences in light-front wave functions and their norm squares.
- 2. Leading divergences resumed into closed equation in planar limits.
- 3. Generates BFKL, BK-like equations in various limits/approximations.



• Emission of the hardest soft gluon splits the original dipole into two dipoles in which softer gluons emit independently.

- Entanglement in Mueller's dipole:
- 1. At order α_s , rapidity divergence leads to enhanced logarithmic behavior: $S \sim \alpha_s \ln^2 N + \alpha_s \ln N$.
- 2. Generally, to order $k: S \sim \alpha_s^k \ln^{k+1} N + \alpha_s^k \ln^k N$.
- 3. Needs re-summation.
- 4. Evolution equation for reduced density matrix can be written out.

$$\begin{split} & \rho_1(b_{10},x_0,x_0',x_{\min},u) = S^{\frac{1}{2}} \bigg(b_{10},\frac{x_0 x_0'}{x_{\min}^2} \bigg) |0\rangle\langle 0| \\ & + \frac{\alpha_s C_F}{\pi^2} \int db_2^2 \frac{b_{10}^2}{b_{12}^2 b_{20}^2} \sum_{x_{\min} < x_1,x_1' < x_0} S^{\frac{1}{2}} \bigg(b_{10},\frac{x_0 x_0'}{x_1 x_1'} \bigg) \frac{|x_1\rangle\langle x_1'|}{\sqrt{x_1 x_1'}} \otimes \rho_1(b_{12},x_1,x_1',x_{\min}) \otimes \rho_1(b_{20},x_1,x_1',x_{\min}) \;. \end{split}$$



• A 1D toy model.

$$Z(y,u) = e^{-ay} + aue^{-ay} \int_0^y e^{ay_1} dy_1 Z(y_1,u) Z(y_1,u) .$$

- 1. $Z(y,u) = \sum_{n=0}^{\infty} u^n p_n(y)$
- 2. Probability of finding n soft gluon: $p_n(y) = e^{-ay}(1 e^{-ay})^n$.
- 3. At large $y \sim \ln x_0 N$, a very wide width $\sim e^{ay}$ of $p_n(y)$.

•
$$p_n(y) = e^{-ay}(1 - e^{-ay})^n$$

- 4. $\langle n \rangle \sim e^{ay}$: exponential in rapidity gap.
- 5. $S(y) y \sim \ln(n) = a \ln x_0 N$. Linear but enhanced.
- 6. $S(y) y = \ln(n)$, the total dipole number $\langle n \rangle$ probed in inclusive process.

• Another interesting example: 1+2D QCD

$$Z(b, y, u) = e^{-mby} + um \int_0^y dy_1 e^{-mb(y-y_1)} \int_0^b db' Z(b - b', y_1, u) Z(b', y_1, u) ,$$

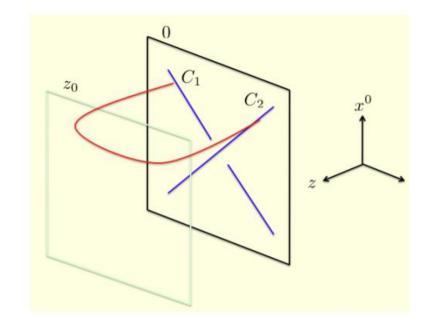
- 1. Phase-space constraint: emitted soft gluon must be within the original dipole!
- 2. Unique to 1+2D. One dimensional transverse direction.

$$\int_{-\infty}^{\infty} \frac{dk^z}{(2\pi)} \frac{e^{ik^z b}}{k^z} = \frac{i}{\pi} \int_0^{\infty} \frac{dk^z}{k^z} \sin(k^z b) = \frac{i}{2} \operatorname{sign}(b) ,$$

- 3. Distribution of soft gluons is Poissonian: $p_n = e^{-mby} \frac{(mby)^n}{n!}$
- Peak at $\langle n \rangle = mby$. Linear instead of exponential.
- Much narrower width: $\langle \delta n^2 \rangle^{\frac{1}{2}} \sim \sqrt{mby}$.
- 4. The entropy $S y = \frac{1}{2} \ln(2\pi e \ mby)$.
- 5. Quenching of phase space, "Kinematic saturation".

Rapidity space entanglement: string picture

- In the string picture, parton-parton scattering depicted by exchange of a minimal surface.
- 1. World-sheet instanton and thermal entropy S_T .
- 2. Quantum entropy $S_E = S_T = \ln \langle N \rangle$.
- 3. $S_E = \frac{1}{6} \frac{\langle \delta b_{\perp}^2 \rangle}{l_S^2}$ like Bekenstein-Hawkins black-hole.

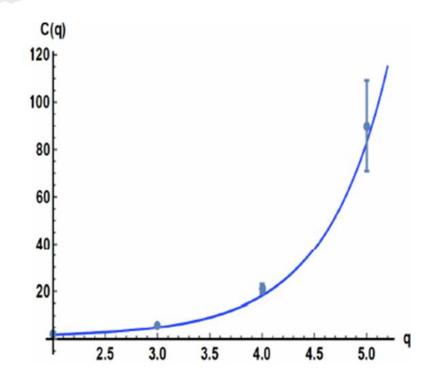


Rapidity space entanglement: string picture

- $\frac{dS_E}{dy} \le \frac{D_\perp}{6}$: chaos bound saturated.
- Cascade equation for particle multiplicities.

$$p_n(D_\perp) = \frac{(n+D_\perp - 1)!}{n!D_\perp!} e^{-\frac{D_\perp}{6}y} \left(1 - e^{-\frac{1}{6}y}\right)^n,$$

$$C(q) = \frac{\langle n^q \rangle}{\langle n \rangle^q} = \frac{1}{\langle n \rangle^q (\langle n \rangle - 1)} \operatorname{PolyLog} \left(-q, 1 - \frac{1}{\langle n \rangle} \right) ,$$



Conclusion

- Rapidity space entanglement as a probe of light-front limit of QFT.
- Subcritical system: S(y) = cy + b with c < 1.
- Critical system: y^n in perturbative expansion. $S(y) y = S_1(y)$ with $S_1(y)$ growth with y.
- 1D reduction: $S_1(y) = ay$. Expected for 4D QCD. Consistent with the string picture.
- Kinematic saturation in 1+2 QCD. $S_1(y) \sim \frac{1}{2} \ln y$.
- Measured by particle multiplicity.